

BULLETIN OF THE CHEMICAL SOCIETY OF JAPAN, VOL. 46, 2227—2229 (1973)

## The CNDO-type Molecular Orbital Calculations of $\text{MnO}_4^-$ , $\text{CrO}_4^{2-}$ , $\text{PdX}_4^{2-}$ , and $\text{PdX}_6^{2-}$

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(Received August 3, 1972)

Since the molecular orbital (MO) calculations of  $\text{MnO}_4^-$  and  $\text{CrO}_4^{2-}$  were presented by Wolfsberg and Hermholz,<sup>1)</sup> various MO methods have been used in the MO calculations of transition-metal complexes; the Extended Hückel,<sup>2)</sup> SCCC,<sup>3)</sup> CNDO-type,<sup>4)</sup> INDO-type,<sup>5)</sup> *ab-initio*,<sup>6)</sup> and the others<sup>7)</sup> methods. However, none of these methods was perfectly satisfactory, because an effective method for some compounds can not always give reasonable results in other cases. In order to re-examine the usefulness of the CNDO-type MO calculation of transition-metal complexes except for the estimation of the intra-atomic triplet transition, we calculated the electronic structures of  $\text{MnO}_4^-$ ,  $\text{CrO}_4^{2-}$ ,  $\text{PdX}_4^{2-}$ , and  $\text{PdX}_6^{2-}$  (X=halogen), for the CNDO-

type approximation avoids the difficulties of the rotation invariance and of the estimation of the one-center exchange integral. The calculated results of the electronic population, the strength of the Pd-X bonds, and the electronic spectra of these complexes will be presented in this paper.

### Method

We start from Roothaan's familiar closed-shell SCF equation<sup>8)</sup> and separate the Hartree-Fock matrix element,  $F_{rs}$ , into the one-electron part,  $H_{rs}$ , and the two-electron part;

$$|F_{rs} - \epsilon S_{rs}| = 0$$

$$F_{rs} = H_{rs} + \sum_{t,u}^{\text{occ}} P_{tu}[(rs|tu) - 0.5(rt|su)] \quad (1)$$

where r, s, t, and u denote the atomic orbitals and whose  $P_{tu}$  is the t, u density matrix element. Then, the following approximations are introduced into Eq. (1): (I) The zero-differential overlap approximation<sup>9)</sup> is introduced:

1) M. Wolfsberg and L. Hermholz, *J. Chem. Phys.*, **20**, 837 (1952).

2) C. J. Ballhausen and H. B. Gray, *Inorg. Chem.*, **1**, 111 (1962).

3) R. F. Fenske and C. C. Sweeney, *ibid.*, **3**, 1105 (1964).

4) a) R. D. Brown, B. H. James, M. F. O'Dwyer, and K. R. Roby, *Chem. Phys. Lett.*, **1**, 459 (1967). b) R. D. Brown, B. H. James, and M. F. O'Dwyer, *Theor. Chim. Acta*, **17**, 262, 279 (1970).

5) a) L. Oleari, G. De Michelis, and L. De Sipio, *Mol. Phys.*, **10**, 111 (1966). b) W. Th. A. M. Van Der Lugt, *Chem. Phys. Lett.*, **10**, 117 (1971).

6) I. H. Hillier and V. R. Saunders, *Proc. Roy. Soc. Ser. A*, **320**, 161 (1970).

7) H. Katô, *This Bulletin*, **45**, 1281 (1972).

8) C. C. J. Roothaan, *Rev. Mod. Phys.*, **23**, 69 (1951).

9) J. A. Pople, D. P. Santry, and G. A. Segal, *J. Chem. Phys.*, **43**, s 129 (1965).

$$S_{rs} = \delta_{rs} \quad (2)$$

$$(rs|tu) = \delta_{rs}\delta_{tu}(rr|tt) \quad (3)$$

(II) The diagonal element,  $H_{rr}$ , of the one-electron part can be written as follows:<sup>10)</sup>

$$H_{rr} = U_{rr} + \sum_{B \neq A} (B|rr) \quad (4)$$

$$U_{rr} = -I_r - (N_r - 1)(rr|rr) - \sum_{s \neq r}^A N_s (rr|ss) \quad (5)$$

$$(B|rr) = - \sum_s^B N_s (rr|ss) \quad (6)$$

where  $N_r$  and  $I_r$  represent the number of electrons on the atomic orbital (AO),  $r$ , in the neutral atom,  $A$ , and the ionization potential of AO  $r$  respectively.  $\sum^A$  indicates the summation on the  $A$  atom.

(III) The off-diagonal element,  $H_{rs}$  ( $r \neq s$ ), is approximated according to Wolfsberg and Hermholz:<sup>11)</sup>

$$H_{rs} = -0.5kS_{rs}(I_r + I_s) \quad (7)$$

where  $k$  is a parameter equal to 1.1.

(IV) The one-center Coulomb repulsion integral  $(rr|rr)$  and the two-center Coulomb repulsion integral  $(rr|ss)$  are estimated by Pariser's approximation<sup>11)</sup> and Ohno's approximation<sup>12)</sup> respectively.

In order to calculate  $S_{rs}$ , Clementi's AO<sup>13)</sup> and Burns' AO<sup>14)</sup> were used for ligands and metals respectively. The same orbital exponents were used for  $s$ - and  $p$ -orbitals. The values obtained by Hinze and Jaffé,<sup>15)</sup> were used for the  $I_r$  and  $(rr|rr)$  values of ligands. The values used for the metals are listed in Table 1.

TABLE 1. THE IONIZATION POTENTIAL ( $I_r$ ) AND ONE-CENTER COULOMB REPULSION INTEGRAL ( $J_{rr}$ )

		$I_r$ (eV)	$J_{rr}$ <sup>18)</sup> (eV)
Cr	3d	-20.28 <sup>16)</sup>	14.03
	4s	-14.73	6.80
	4p	-10.31	5.80
Mn	3d	-21.64 <sup>16)</sup>	14.70
	4s	-15.35	7.09
	4p	-10.62	5.95
Pd	4d	-8.54 <sup>17)</sup>	13.87
	5s	-7.60	7.46
	5p	-3.84	6.22

10) T. Yonezawa, K. Yamaguchi, and H. Kato, *This Bulletin*, **40**, 536 (1967); H. Kato, H. Konishi, H. Yamabe, and T. Yonezawa, *ibid.*, **40**, 2761 (1961).

11) R. Pariser, *J. Chem. Phys.*, **21**, 568 (1953).

12) K. Ohno, *Theor. Chim. Acta*, **2**, 219 (1964).

13) E. Clementi and D. L. Raimondi, *J. Chem. Phys.*, **38**, 2686 (1963).

14) G. Burns, *ibid.*, **41**, 1521 (1964).

15) J. Hinze and H. H. Jaffé, *J. Amer. Chem. Soc.*, **84**, 540 (1961); *J. Phys. Chem.*, **67**, 1504 (1963).

16) C. J. Ballhausen and H. B. Gray, "Molecular Orbital Theory," W. A. Benjamin, Inc., New York (1965), p. 120.

The electron configuration of neutral Mn atom is  $d^5s^2$ . However, we can not obtain the  $I_r$  for  $d^5s^2$  except with  $I_{4s}$ , because of the lack of the observed atomic spectra. Therefore we investigated the effect of  $I_r$  and choose the suitable values. We show this value in Table 1.

17) I. H. Hillier and R. M. Canadine, *Discuss. Faraday Soc.*, **47**, 27 (1969).

18) L. De Sipio, E. Tondello, G. De Michelis, and L. Oleari, *Chem. Phys. Lett.*, **11**, 287 (1971).

## Results and Discussion

**$MnO_4^{2-}$  and  $CrO_4^{2-}$ .** For the purpose of comparing our results with previous ones,<sup>4-6)</sup> we calculated the MO's of  $MnO_4^{2-}$  and  $CrO_4^{2-}$ . The net charge on the Mn atom in  $MnO_4^{2-}$  was found to be 1.28 from the metal K-X ray spectra.<sup>4b)</sup> However, in our results, negative net charge on the Mn atom was obtained. This may be due to the too large electron population of the 4s and 4p orbitals of the Mn atom compared with the results of Oleari *et al.*<sup>5a)</sup> and Hillier and Saunders<sup>6)</sup> ( $MnO_4^{2-}$ ; 4s, 0.81 and 4p, 0.78;  $CrO_4^{2-}$ ; 4s, 0.81 and 4p, 0.78 in our results).

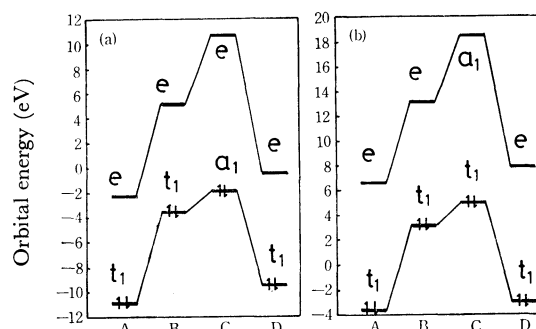


Fig. 1. HOMO and LUMO of  $MnO_4^{2-}$ (a) and  $CrO_4^{2-}$ (b) by various Methods.

||: HOMO, —: LUMO

A: Oleari *et al.*<sup>5b)</sup> by INDO, B: Brown *et al.*<sup>4b)</sup> by VE-SCF-CNDO, C: Hillier *et al.*<sup>6)</sup> by *ab-initio*, D: ours by CNDO.

The highest occupied (HO) MO and the lowest unoccupied (LU) MO are considered to be represented by  $t_1^{4,5)}$  and  $e^{19)}$  respectively. Our results with respect to HOMO and LUMO are compared with previous ones in Fig. 1. The results of Oleari *et al.*<sup>5a)</sup> seem to be quite reasonable. However, the calculation by Brown *et al.*<sup>4b)</sup> gave a positive orbital energy for the HOMO of  $CrO_4^{2-}$ . Also, in the case of Hillier's calculation,<sup>6)</sup> the representations of HOMO and LUMO are not  $t_1$  and  $e$  respectively, and the HOMO of  $CrO_4^{2-}$  has a positive orbital energy. In contrast to these results, our results are reasonable regarding the orbital energy and the orbital symmetry.

The calculated transition energies are given in Table 2. Our results agree fairly with the experimental results.

**$PdX_4^{2-}$  and  $PdX_6^{2-}$ .** We then calculated the MO's of  $PdX_4^{2-}$  and  $PdX_6^{2-}$  in order to examine the validity of our method for different types of complexes. The net charge of Pd decreases with a decrease in the electronegativity of halogens, as is shown in Table 3. This tendency agrees with the results obtained by means of photoelectron spectroscopy.<sup>21)</sup> However, we obtained a negative net charge on the Pd atom due to the large electron population of 5s- and 5p-orbitals. As is shown in Table 3, the absolute value of the two atomic part of the total energy,  $|E_{Pd-Cl}|$ ,<sup>9)</sup> is larger than  $|E_{Pd-Br}|$  in  $PdX_4^{2-}$  and  $PdX_6^{2-}$ , and  $|E_{Pd-X}|$  in

19) A. Carrington, D. J. E. Ingram, K. A. K. Lott, D. S. Schonland, and M. C. R. Symons, *Proc. Roy. Soc., Ser. A*, **254**, 101 (1960).

TABLE 2. TRANSITION ENERGIES OF  $\text{MnO}_4^{2-}$ ,  $\text{CrO}_4^{2-}$ , AND  $\text{PdCl}_6^{2-}$ 

	Experimental (eV)	Calculated (eV)
$\text{MnO}_4^{2-}$ <sup>20b)</sup>	1.79 ${}^3\text{T}_1$ or ${}^1\text{T}_1$ <sup>20a)</sup>	1.44 ${}^1\text{T}_1$
	2.38 ${}^1\text{T}_2$	2.45 ${}^1\text{T}_2 + {}^1\text{E}$
		3.65 ${}^1\text{T}_2$
	3.44 ${}^1\text{T}_1$	4.27 ${}^1\text{T}_1$
$\text{CrO}_4^{2-}$ <sup>20c)</sup>	4.13 ${}^1\text{T}_2$	4.47 ${}^1\text{T}_2$
	3.00 ${}^1\text{T}_1$	3.36 ${}^1\text{T}_1$
		3.78 ${}^1\text{T}_1$
	3.56 ${}^1\text{T}_2$	5.15 ${}^1\text{T}_2$
$\text{PdCl}_6^{2-}$ <sup>24)</sup>	4.65 ${}^1\text{T}_2$	6.47 ${}^1\text{T}_2$
	6.20	6.70 ${}^1\text{T}_1$
	2.27	3.98 $\text{T}_{1g}$
	3.64	4.13 $\text{T}_{2g}$
		5.94 $\text{T}_{1u}$
	5.17	6.45 $\text{T}_{2u}$

$\text{PdX}_6^{2-}$  is larger than  $|E_{\text{Pd-X}}|$  in  $\text{PdX}_4^{2-}$ . These results agree with the experimental results.<sup>22)</sup>

Though the order of bond strength agrees with the experimental results, we failed to obtain reasonable transition energies for  $\text{PdX}_4^{2-}$ . The calculated first transition energies seem too small for  $\text{PdCl}_4^{2-}$  (0.019 eV) and  $\text{PdBr}_4^{2-}$  (0.028 eV). Furthermore, neither is the d-d transition which is to be expected as the first transition.<sup>23)</sup> For  $\text{PdCl}_6^{2-}$ , the calculation gives almost reasonable results, as is shown in Table 2.

As has been described above, by treatment with the same approximation as was used in  $\text{MnO}_4^{2-}$  and  $\text{CrO}_4^{2-}$  we can obtain reasonable trends of electron population and bond strength for Pd-halogen complexes, but not successful results for any transition energies except  $\text{PdX}_6^{2-}$ . Attempts to improve this approximation are under way, for example the inclusion of one-center exchange integrals and further refinements in the parametrization of  $H_{rs}$  will be described in the near future.

TABLE 3. ELECTRON POPULATION AND  $E_{\text{Pd-X}}$  OF  $\text{PdX}_4^{2-}$  AND  $\text{PdX}_6^{2-}$ 

	Electron Population			Net Charge on Pd		$E_{\text{Pd-X}}^{c)}$	Obsd <sup>22)</sup>	
	Pd		X	Calcd	Expt. <sup>21)</sup>		f <sub>Pd-X</sub> <sup>a)</sup>	
							ν <sub>Pd-X</sub> <sup>b)</sup>	
	d	sp						
PdCl <sub>4</sub> <sup>2-</sup>	8.075	2.259	7.417	−0.333	0.42	−15.882 <sub>ev</sub>	1.355	310
PdBr <sub>4</sub> <sup>2-</sup>	8.049	2.343	7.402	−0.392	0.23	−14.688		187
PdF <sub>6</sub> <sup>2-</sup>	7.137	2.219	7.441	0.644		−19.124		
PdCl <sub>6</sub> <sup>2-</sup>	7.177	3.043	7.297	−0.220	0.72	−16.300	1.90	317
PdBr <sub>6</sub> <sup>2-</sup>	7.157	3.203	7.273	−0.360		−14.920	1.58	198

a) The force constant of Pd-X stretching vibration. (mdyn/Å)

b) The symmetric stretching frequency of Pd-X. ( $\text{cm}^{-1}$ )

c)  $E_{A-B} = 2 \sum_r^A \sum_s^B P_{rs} H_{rs} - \frac{1}{2} \sum_r^A \sum_s^B P_{rs}^2 (\text{rr/ss}) + \sum_r^A \sum_s^B (P_{rr} - N_r) (P_{ss} - N_s) (\text{rr/ss})$

20) a) P. Day, L. De Sipio, and L. Oleari, *Chem. Phys. Lett.*, **5**, 533 (1970); b) L. W. Johnson and S. P. McGlynn, *J. Chem. Phys.*, **55**, 2985 (1971); c) L. W. Johnson and S. P. McGlynn, *Chem. Phys. Lett.*, **7**, 618 (1970).

21) G. Kumer, J. R. Blackburn, R. G. Albridge, W. E. Modderman, and M. M. Jones, *Inorg. Chem.*, **11**, 296 (1972).

22) P. J. Hendra and P. J. D. Park, *Spectrochim. Acta*, **23A**, 1635 (1967).

23) A. J. MacCaffery, P. N. Schatz, and P. J. Stephens, *J. Amer. Chem. Soc.*, **90**, 5730 (1968).

24) C. K. Jorgensen, *Adv. Chem. Phys.*, **5**, 119 (1963); *Mol. Phys.*, **2**, 309 (1959).